

## 195. The Photoelectron Spectra of 4, 13-Diaza- and 4, 16-Diaza [2<sub>4</sub>](1, 2, 4, 5)cyclophanes<sup>1)</sup>

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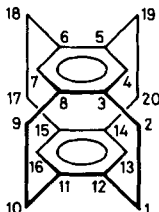
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### Summary

A tentative assignment is proposed for the He(I $\alpha$ ) photoelectron (PE.) spectra of the title compounds. It is based on the assumption that the first four  $\pi$ -bands occupy the same position in both spectra, as suggested by a qualitative correlation with the PE. spectrum of [2<sub>4</sub>](1, 2, 4, 5)cyclophane and by a previously proposed molecular orbital model. If this is accepted, then the bands due to electron ejection from the lone-pair orbitals are split by roughly 0.8 eV in the case of the 4, 13-diaza isomer, whereas no split can be detected in the case of the 4, 16-diaza isomer.

**1. Introduction.** – The PE. spectra of pyridine [1–3] and of benzene [2] differ in that the  $e_{1g}^{-1}$  band of the latter (at 9.25 eV) is split in the spectrum of pyridine into two bands at  $\sim 9.7$  and 10.5 eV, corresponding to  $a_2(\pi)^{-1}$  and  $b_1(\pi)^{-1}$  ionization processes, respectively. In addition, there is a band at  $\sim 9.7$  eV due to the removal of an electron from the lone-pair orbital  $a_1(n)$ . In this note we investigate,

- 1) The nomenclature refers to the following convention for the numbering of the centres in [2<sub>n</sub>]cyclophanes, in particular in [2<sub>4</sub>](1, 2, 4, 5)cyclophane:

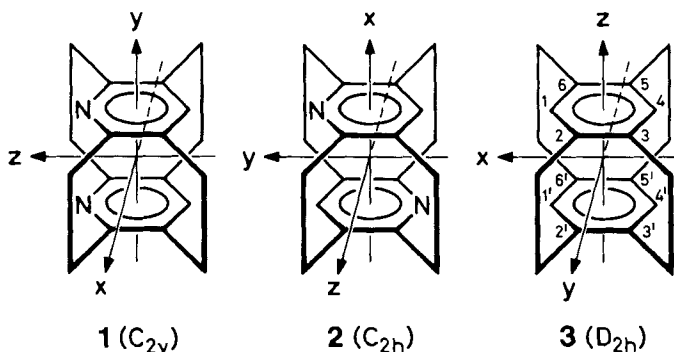


However, for the sake of convenience and simplicity we shall depart from this convention within the text of the present contribution.

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how the above ionization energies are affected when two pyridine nuclei are held closely together in a cyclophane system. This should provide ancillary information concerning previous assignments of cyclophane PE. spectra [4-6].

**2. Experimental Results.** - In *Figure 1* are presented the He(I $\alpha$ ) PE. spectra of 4,13-diaza[2<sub>4</sub>](1,2,4,5)cyclophane (**1**), 4,16-diaza[2<sub>4</sub>](1,2,4,5)cyclophane (**2**) [7] and of [2<sub>4</sub>](1,2,4,5)cyclophane (**3**) [4]. To simplify the ensuing discussion, we re-number the centres of the upper (u) and lower (l) deck in **3**, and thus in **1** and **2**, as indicated below, so that corresponding  $\pi$ -centres carry the same numbers  $\mu, \mu'$ .



In the ionization energy interval  $8 \text{ eV} < I < 10 \text{ eV}$ , the PE. spectrum of the *syn*-isomer **1** consists of three maxima with integrated intensity ratios of 3:2:1, whereas the *anti*-isomer **2** shows four maxima with ratios 1:3:1:1. The positions of these maxima and thus, approximately, the ionization energies  $I_j^m$  corresponding to the underlying individual bands, are given in *Table 1*, together with the relevant data for the parent hydrocarbon **3** [4] and with the tentative assignment to be discussed below.

**3. Discussion.** - Clearly the problem of assigning the bands in the interval  $8 \text{ eV} < I < 10 \text{ eV}$  is made difficult by their strong overlap and, in fact, the following analysis leads to a sequence of the radical cation states for **1**<sup>+</sup> and **2**<sup>+</sup> which should be regarded, at best, as a plausible working hypothesis.

We begin by discussing those orbitals of **1** and **2** which consist essentially of the linear combinations of the semi-localized  $\pi$ -orbitals of the two pyridine rings. In the following we designate the  $\pi$ -orbitals  $a_2(\pi)$  and  $b_1(\pi)$  of the upper and lower pyridine nuclei by  $S_u, S_l$  and  $A_u, A_l$  respectively, these orbitals being symmetric (S) or antisymmetric (A) relative to a mirror-plane containing the centres 1, 1' and 4, 4'. (The analogue nomenclature has been used in the previous contributions [4] [5], to which the reader is referred for further details.) From these we form symmetry adapted linear combinations  $S_{\pm} = (S_u \pm S_l) / \sqrt{2}$  and  $A_{\pm} = (A_u \pm A_l) / \sqrt{2}$ . In turn, these are the major contributors to the following molecular orbitals of **1** and **2**, the symmetry labels of which refer to the coordinate systems included in formulae 1-3 (see *Table 2*).

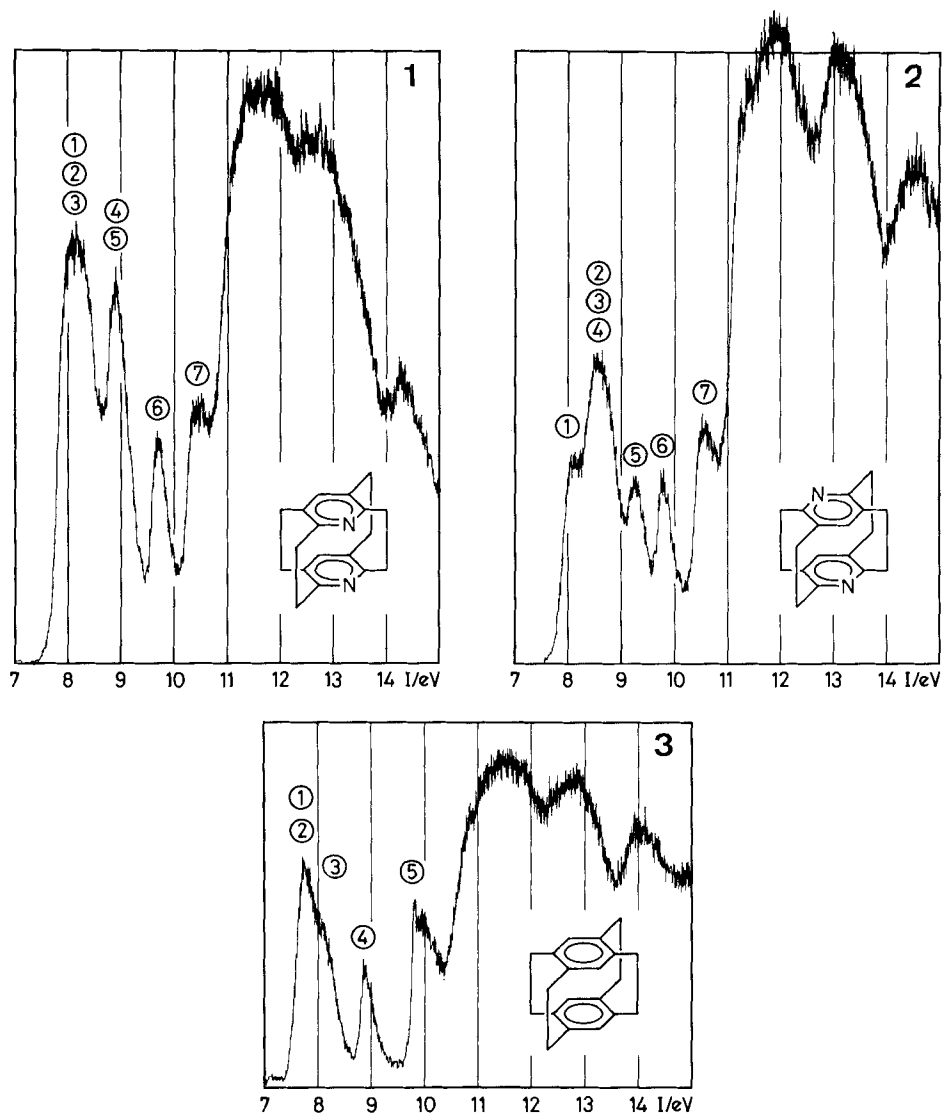


Fig. 1. *He(I $\alpha$ ) PE spectra of 4,13-diaza[24](1,2,4,5)cyclophane (1), 4,16-diaza[24](1,2,4,5)cyclophane (2) and of [24](1,2,4,5)cyclophane (3)*

For comparison, the corresponding orbitals of the parent hydrocarbon **3** have been included, in which case  $A_{\pm}$  and  $S_{\pm}$  refer to the linear combinations of the benzene  $\pi$ -orbitals. Note that  $A_{+}$  and  $S_{+}$  are antisymmetric,  $A_{-}$  and  $S_{-}$  symmetric, with respect to a reflection in the plane passing between the two decks.

Two ways present themselves for the assignment of orbital energies to the orbitals of **1** and **2** listed in *Table 2*. The first one is entirely empirical and based

Table 1. Tentative assignment and ionization energies of the PE. spectra of 1, 2 and 3

Band	1		2		3 [4]	
	Orbital	$I^m/\text{eV}$	Orbital	$I^m/\text{eV}$	Orbital	$I^m/\text{eV}$
①	$a_2(\pi)$	8.2	$b_g(\pi)$	8.1	$b_{3g}(\pi)$	7.7
②	$b_1(\pi); b_2(n)$		$a_u(\pi)$	8.6	$b_{2g}(\pi)$	
③			$a_g(n); b_u(n)$		$b_{2u}(\pi)$	8.1
④	$a_1(n)$	8.9			$b_{3u}(\pi)$	8.8
⑤	$b_2(\pi)$		$a_g(\pi)$	9.3		9.8
⑥	$a_1(\pi)$	9.7	$b_u(\pi)$	9.8		
⑦		10.5		10.5		

on a first-order perturbation argument, starting with the orbitals of 3 [4] (*cf.* Table 1). Replacement of a CH-unit in benzene ( $I(e_{1g}^{-1})=9.25$  eV) by N, to yield pyridine ( $I(a_2^{-1})=9.7$  eV;  $I(b_1^{-1})=10.5$  eV) shifts the orbital A by  $-0.45$  eV and orbital S by  $-1.25$  eV [1-3]. Similarly, the A and S orbitals of *m*-xylene ( $I(a_2^{-1})=8.5$  eV;  $I(b_1^{-1})=9.0$  eV [8]) are shifted by  $-0.4$  and  $-1.0$  eV respectively, if the CH-group in position 2 is replaced by N to yield 2,6-dimethylpyridine ( $I(a_2^{-1})=8.9$  eV;  $I(b_1^{-1})=10.0$  eV [3]). Assuming in a crude approximation that similar shifts affect the  $\pi$ -orbital energies of 3 when two N-atoms are introduced to yield 1 or 2, then the orbital energies given by the expressions 1 are approximately expected for the latter molecules.

$$\begin{array}{l}
 \mathbf{1} \quad \mathbf{2} \quad \mathbf{3} \\
 \varepsilon(a_2) = \varepsilon(b_g) = \varepsilon(b_{3g}) - 0.4 \text{ eV} = -8.1 \text{ eV} \\
 \varepsilon(b_1) = \varepsilon(a_u) = \varepsilon(b_{2u}) - 0.4 \text{ eV} = -8.5 \text{ eV} \\
 \varepsilon(b_2) = \varepsilon(a_g) = \varepsilon(b_{2g}) - 1.2 \text{ eV} = -8.9 \text{ eV} \\
 \varepsilon(a_1) = \varepsilon(b_u) = \varepsilon(b_{3u}) - 1.2 \text{ eV} = -10.0 \text{ eV}
 \end{array} \quad (1)$$

The major shortcomings of this naive procedure, which necessarily predicts the same energies for corresponding orbitals of 1 and 2, are that it assumes the same geometry for all three compounds and that it neglects second-order terms which one might feel to be important. To take these factors into account we performed model calculations along the lines described previously [4-6].

To this end we have first parametrized a *Hückel*-type model for pyridine, to reproduce its observed  $\pi^{-1}$ -ionization energies. This yields *ad hoc* the atomic Coulomb integrals  $a_N = -9$  eV,  $a_{C(\text{ortho to N})} = -6.7$  eV and the resonance integral  $\beta_{CN} = -2.5$  eV, if all other values are kept unchanged, as described previously [5]:  $a_C = -5.8$  eV;  $\beta_{C,C} = -3.2$  eV;  $a_{\text{CH}_2\text{CH}_2\text{-bridge}} = -14.0$  eV;  $\beta_{C,\text{CH}_2\text{CH}_2\text{-bridge}} = -2.4$  eV.

Table 2

Type	1 ( $C_{2v}$ )	2 ( $C_{2h}$ )	3 ( $D_{2h}$ )
A <sub>+</sub>	a <sub>2</sub>	b <sub>g</sub>	b <sub>3g</sub>
A <sub>-</sub>	b <sub>1</sub>	a <sub>u</sub>	b <sub>2u</sub>
S <sub>+</sub>	b <sub>2</sub>	a <sub>g</sub>	b <sub>2g</sub>
S <sub>-</sub>	a <sub>1</sub>	b <sub>u</sub>	b <sub>3u</sub>

The crystal structure analysis of the compounds **1** [9], **2** [10], **3** [11] reveals that in all three compounds the centres 1, 1' and 4, 4' are further apart than the bridged ones: **1**: 1, 1', 284 pm; 4, 4', 294 pm; mean of others 265 pm; **2**: 1, 1' and 4, 4', 287 pm; others 264 pm; **3**: 1, 1', 295 pm, 2, 2', 269 pm. Using these experimental geometries and the parameters given above leads again to the result that the positions of the first four  $\pi$ -bands in the PE. spectra of **1** and **2** should be the same within rather narrow limits of error (see Table 3).

Table 3

	Orbital $\epsilon/\text{eV}$	a <sub>2</sub> -8.28	b <sub>1</sub> -8.57	b <sub>2</sub> -8.66	a <sub>1</sub> -10.29
<b>1</b>					
<b>2</b>	Orbital $\epsilon/\text{eV}$	b <sub>g</sub> -8.27	a <sub>u</sub> -8.57	a <sub>g</sub> -8.71	b <sub>u</sub> -10.03

In other words, the relative orientation of the two pyridine rings in the molecules **1** and **2** is predicted to have only a negligible influence on the orbital energies. The same phenomenon has already been observed in the case of the *syn*- and *anti*-[2.2](1,6)azulenophanes [6]. This makes the naive perturbation model underlying the results 1 more palatable and it is indeed surprising that both treatments yield the same sequence of orbitals and rather close-lying orbital energies.

Taking the above results at face value and postulating that the first four  $\pi$ -bands occupy the same positions in the PE. spectra of **1** and **2**, the following procedure suggests itself. We first normalize the spectra of **1** and **2** in the interval  $7 \text{ eV} < I < 10.5 \text{ eV}$  in such a way that the integrated intensity is the same for both. This is equivalent to the assumption that the intensities of corresponding pairs of the six bands ① to ⑥ is the same (*cf.* Fig. 2). Then the two spectra are shifted by 0.1 eV relative to each other, matching the onsets, and then subtracted, *i.e.* the spectrum of **2** from that of **1**. In the difference curve labeled **1-2** in Figure 2, the positive maxima should then correspond to the lone-pair bands in the PE. spectrum of **1**, the negative minimum to the corresponding bands in **2**. For rather obvious reasons such a treatment is very crude, but the result seems pleasingly clear-cut. The difference curve suggests that in **1** the two lone-pair orbitals, *i.e.* those molecular orbitals which are dominated by the linear combinations  $n_{\pm} = (n_1 \pm n_1')/\sqrt{2}$ , are split by  $\sim 0.8 \text{ eV}$ , whereas the corresponding orbitals in **2**, based on  $n_{\pm} = (n_1 \pm n_4')/\sqrt{2}$  exhibit the same orbital energy within narrow but unspecified limits. In contrast to what the smoothed curves of Figure 2 suggest, the procedure

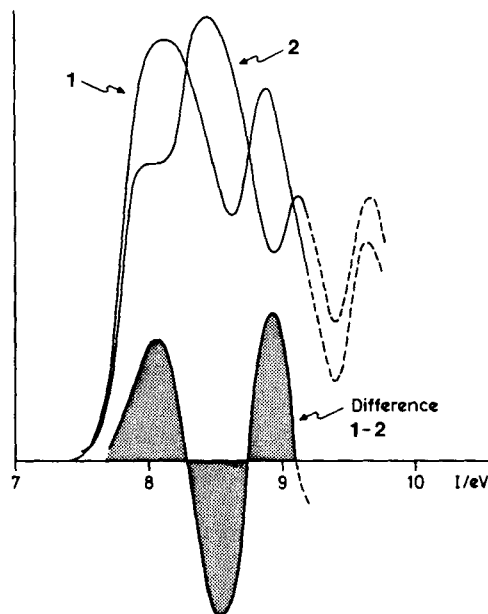


Fig. 2. Normalized PE. spectra of **1** and **2** and the difference spectrum **1-2**

outlined above depends critically on the numerous assumptions implied and not too much significance can be attached to the positions of the maxima (8.1 eV; 8.9 eV) and the minimum (8.5 eV) of the difference curve **1-2**. Notwithstanding the rather large uncertainties, the values are of the size expected, if compared to the positions of the lone-pair bands in the PE. spectra of alkyl-substituted pyridines [1][3].

Intuitively it seems reasonable that the two lone-pair orbitals  $n_1$ ,  $n_4'$  do not interact appreciably in **2**, whereas an interaction, yielding a split of  $\sim 0.8$  eV exists between  $n_1$  and  $n_1'$  in **1**, where the two lone-pairs face each other across a distance of 284 pm [9]. Interactions of this size between lone-pair orbitals have been observed in [2.2](2,6)metapyridinophane [12] and are also suggested by the PE. spectra of other heterocyclic [2.2]cyclophanes [13]. However, the origin of a split of  $\sim 0.8$  eV is by no means obvious. From the computed overlap  $S_{nn} = \langle n_1 | n_1' \rangle \approx 0.01$  to 0.02, depending on the assumptions made, one would expect a 'through-space' contribution of 0.2 to 0.4 eV towards the observed split, *i.e.* about 1/4 to 1/2 of its size. The difference must then be assigned to a 'through-bond' interaction involving five intermediate C,C-bonds on each side of the pair  $n_1$ ,  $n_1'$ . Once more, the magnitude of such a contribution is in keeping with expectation based on previous experience [14].

The resulting, tentative assignment is summarized in *Table 1*. It is obvious from the preceding discussion that it must be taken with more than only one grain of salt and that it should be accepted, at best, as a reasonable working hypothesis, which we hope is not too far from the truth.

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